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Gyrotactic trapping of micro-swimmers in simple shear flows: A study directly from the fundamental Smoluchowski equation

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Objectives

Thin phytoplankton layers with vertically compressed structures found in stratified lakes and coastal water

- Deriving the eigenfunction and dual basis-functions by a Galerkin method
- Biorthogonal expansion of moments

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bodies have been a hot-spot in marine science and fluid mechanics. Although extensive efforts have been made to explore gyrotactic trapping as a possible mechanism, obvious inconsistencies remain there between the generalized Taylor dispersion method and individualbased model. In this work, a study directly from the fundamental Smoluchowski equation is carried out on the gyrotactic trapping mechanism in representative simple shear flows.

Methods



The problem setting

• Governing equation (Smoluchowski equation)

$$\frac{\partial P}{\partial t} = -[Pe_s \sin \phi \cos \theta + Pe_f U(z)] \frac{\partial P}{\partial x} + D_t \frac{\partial^2 P}{\partial x^2} + \mathcal{L}_u P,$$

• *Reflective boundary conditions*

Results



Figure 3. Transient vertical distribution of a vertically uniform line-source release of (a) nonsettling (b) settling gyrotactic micro-swimmers in the plane Poiseuille flow.



Figure 4. Transient two-dimensional concentration distributions of a vertically uniform line-source release of non-settling gyrotactic micro-swimmers in a strong plane Poiseuille flow. Left: Edgeworth expansion; right: individual-based model (IBM). Top: t = 5; bottom: t = 10.

$$-2Pe_dP|_{\theta=\theta_0} - D_t \frac{\partial P}{\partial z}\Big|_{\theta=\theta_0} - D_t \frac{\partial P}{\partial z}\Big|_{\theta=2\pi-\theta_0} = 0 \quad \text{on } z = 0, 1, \quad \forall \theta_0 \in [0, 2\pi).$$

Solution procedure

• Transformation of the p.d.f.

$$P(x, z, \phi, \theta, t) = f_r(z)G^r(x, z, \phi, \theta, t),$$
$$f_r(z) = \exp\left(-\frac{Pe_d}{D_t}z\right).$$

• Classical B.C. for the transformed p.d.f.

$$\left. \begin{array}{l} G^{r}|_{\theta} = G^{r}|_{2\pi-\theta}, \\ \frac{\partial G^{r}}{\partial z}|_{\theta} = -\frac{\partial G^{r}}{\partial z}|_{2\pi-\theta}, \end{array} \right\} \quad \text{on} \quad z = 0, 1, \quad \forall \, \theta \in [0, 2\pi)$$

• **Basis functions** $N_n \sqrt{\frac{2l+1}{4\pi}} \cos(n\pi z) P_l(\cos\phi),$ $N_n \sqrt{\frac{2l+1}{2\pi}} \frac{(l-m)!}{(l+m)!} \cos(n\pi z) \cos(m\theta) P_l^m(\cos\phi),$ $N_n \sqrt{\frac{2l+1}{2\pi}} \frac{(l-m)!}{(l+m)!} \sin(n\pi z) \sin(m\theta) P_l^m(\cos\phi),$



Figure 2 Illustration of the reflective boundary conditions.



Figure 5. Comparison of vertical distributions of a vertically uniform line-source release of non-settling gyrotactic micro-swimmers between the current model (BE), Bearon et al. (2011) and IBM.

Conclusions

We have given a more complete picture of gyrotactic trapping. Our results are in good agreement with the IBM simulations due to no approximation made to the fundamental Smoluchowski equation. We found that a steady thin layer can be realized in the solution by adding a small settling speed to counteract the gravitactic focusing at the upper surface.